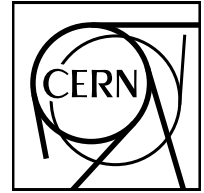


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Measurement of event shapes at large momentum transfer with the ATLAS detector in pp collisions at $\sqrt{s} = 7$ TeV

The ATLAS Collaboration

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The ATLAS Collaboration^{a,1}

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Keywords event shapes · jets · LHC · ATLAS

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1 Introduction

Event shapes represent a generic class of observables that describe the patterns, correlations, and origins of the energy flow in an interaction. In terms of hadronic jet production, event shapes are an indirect probe of multi-jet topologies. These observables have had a long and fruitful history, having been used to measure the strong coupling constant α_s and to test asymptotic freedom [1–4], to constrain color factors for quark and gluon couplings [5], to assess the accuracy of leading order (LO) and next-to-leading order (NLO) Monte Carlo (MC) generators [6, 7], to determine the contribution of non-perturbative quantum chromodynamics (QCD) power corrections [8], and to search for

physics beyond the Standard Model [9]. Furthermore, recent efforts to provide advanced, high-precision theoretical calculations of a range of event shapes for the Large Hadron Collider [10, 11] provide renewed impetus for making such measurements.

This analysis considers six event shapes calculated using hadronic jets. These observables are crucially tied to both the multi-jet nature of the final state produced in high energy collisions and have a strong history in the literature: the third-jet resolution parameter [4, 12–14], y_{23} ; the sphericity and transverse sphericity [15, 16], S and S_\perp ; the aplanarity, A ; and the event thrust and its minor component [17], τ_\perp and $T_{m,\perp}$. Events with high transverse momentum central leading-jet pairs are used for the measurements. Each event shape variable is defined such that it vanishes in the limit of a pure $2 \rightarrow 2$ process and increases to a maximum for uniformly distributed energy within a multi-jet event. Hard gluon emission is thereby signified by large non-zero values of each observable. Furthermore, some of the event shape variables are evaluated as ratios of final state observables, which reduces their sensitivity to jet energy scale (JES) calibration uncertainties as well as other experimental and theoretical uncertainties. These measurements permit detailed tests of the phenomenological models of QCD in leading order MC programs and to indirectly test the running of α_s through measurements performed as a function of the average leading jet momentum. In addition, these results may be used to provide input to tune MC generators in the future.

All event shapes measured in this analysis are defined using jets to represent the final state four-momenta, as discussed in Section 2. The ATLAS detector is described in Section 3, with a particular emphasis on the components relevant for the measurement of event shape variables. Section 4 presents the event selection and description of simulated events which are compared to the data. Jet definitions, calibrations, and selection criteria are also described in Sec-

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tion 4. Finally, the results of these measurements are presented in Section 5.

2 Event Shape Definitions

Six event shapes are measured using high transverse momentum (p_T) jets. The first observable, y_{23} , is a measure of the third-jet p_T relative to the summed transverse momenta of the two leading jets in a multi-jet event and is defined as:

$$y_{23} = \frac{p_{T,3}^2}{H_{T,2}^2}, \quad (1)$$

where $H_{T,2} = (p_{T,1} + p_{T,2})$ is the scalar sum of jet momenta and the subscript $i = 1, 2, 3$ refers to the leading, subleading, or third leading jet in the event. The range of allowed values for y_{23} is $0 \leq y_{23} < 1/4$ and it is often expressed as $\ln y_{23}$ [11, 16]. This definition is different from the original [12] definition which uses the JADE jet algorithm [18]. Eq. 1 is defined with an explicit third-jet as opposed to a continuously variable threshold in the jet algorithm. The sphericity, S , transverse sphericity, S_\perp , and aplanarity, A , embody more global information about the full momentum tensor of the event, M_{xyz} , via its eigenvalues λ_1 , λ_2 and λ_3 :

$$M_{xyz} = \sum_i \begin{pmatrix} p_{xi}^2 & p_{xi}p_{yi} & p_{xi}p_{zi} \\ p_{yi}p_{xi} & p_{yi}^2 & p_{yi}p_{zi} \\ p_{zi}p_{xi} & p_{zi}p_{yi} & p_{zi}^2 \end{pmatrix}, \quad (2)$$

where the sum runs over all jets used in the measurement. The individual eigenvalues are normalized and ordered such that $\lambda_1 > \lambda_2 > \lambda_3$ and $\sum_i \lambda_i = 1$ by definition. These terms are used to define the three observables as

$$S = \frac{3}{2}(\lambda_2 + \lambda_3), \quad (3)$$

$$S_\perp = \frac{2\lambda_2}{\lambda_1 + \lambda_2}, \quad (4)$$

$$A = \frac{3}{2}\lambda_3. \quad (5)$$

Sphericity, Eq. (3), and transverse sphericity, Eq. (4), measure the total transverse momentum with respect to the sphericity axis defined by the four-momenta used for the event shape measurement (specifically, the first eigenvector). The allowed range of S values is $0 \leq S < 1$, but due to the inclusion of the smallest eigenvalue, λ_3 , the typical maximum achieved experimentally is $S \sim 0.8$. Conversely, the transverse sphericity is constructed using the two largest eigenvalues, and the typical range coincides with the allowed range, $0 \leq S_\perp < 1$. Aplanarity (Eq. 5) measures the amount of transverse momentum in or out of the plane formed by the two leading jets via only the smallest eigenvalue of M_{xyz} , λ_3 , with allowed values $0 \leq A < 1/2$. Typical measured values lie between $0 \leq A < 0.3$, with values near zero indicating relatively planar events. The transverse thrust, T_\perp , and its minor component, $T_{m,\perp}$, define a so-called thrust axis for the event, with

respect to which, the total transverse momentum of the jets used in the measurement is minimized. These quantities are defined as

$$T_\perp = \max_{\hat{n}_\perp} \frac{\sum_i |\mathbf{p}_{Ti} \cdot \hat{n}_\perp|}{\sum_i p_{Ti}}, \quad (6)$$

$$\tau_\perp = 1 - T_\perp, \quad (7)$$

$$T_{m,\perp} = \frac{\sum_i |\mathbf{p}_{Ti} \times \hat{n}_\perp|}{\sum_i p_{Ti}}, \quad (8)$$

where T_\perp is translated into τ_\perp in order to maintain a common event shape definition in which a large value indicates a departure from a two-body system. The unit vector \hat{n}_\perp defines the thrust axis of the event. The so-called event plane is defined by \hat{n}_\perp and the beam direction and allows a measurement of $T_{m,\perp}$. The variable $T_{m,\perp}$ quantifies the sum of all transverse momenta \mathbf{p}_{Ti} out of the event plane, where the sum again runs over each jet i considered in the final state. The allowed values for τ_\perp span the range $0 \leq \tau_\perp < 1/3$ due to the range over which both T_\perp and $T_{m,\perp}$ may fall, $0 \leq T_\perp, T_{m,\perp} < 2/3$.

Event shapes constructed using hadronic jets in this way offer several advantages over explicit cross-section calculations for inclusive and multi-jet production. Event shapes may be defined as normalized ratios of hadronic final state observables, thus reducing the sensitivity to experimental uncertainties. Various choices of event shape quantities can also lead to enhanced or suppressed sensitivity to different components of the fundamental physical processes involved [11]. The effect of the underlying event and parton shower can be reduced by focusing only on the leading jets. The choice of renormalization and factorization scales used in calculating the LO and NLO cross-sections may be less important when considering ratios of quantities. Systematic uncertainties, such as the jet energy scale and detector effects, are mitigated by examining the normalized shapes as opposed to absolute cross-sections.

3 The ATLAS Detector

The ATLAS detector [19, 20] provides nearly full solid angle coverage around the collision point¹ with an inner tracking system covering $|\eta| < 2.5$, electromagnetic and hadronic calorimeters covering $|\eta| < 4.9$, and a muon spectrometer covering $|\eta| < 2.7$. Of the multiple ATLAS subsystems, the most relevant to this analysis are the inner tracking detector (ID) [21], the barrel and end-cap calorimeters [22, 23] and the trigger [24].

¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

The ID is comprised of a pixel tracker closest to the beamline, a microstrip silicon tracker, and lastly a straw-tube transition radiation tracker at the largest radii. These systems are layered radially upon each other in the central region. A thin solenoid surrounding the tracker provides an axial 2T field enabling measurement of charged particle momenta.

The calorimeter is built of multiple sub-detectors with several different designs spanning the pseudorapidity range up to $|\eta| < 4.9$. The measurements of event shapes are predominantly performed using data from the central calorimeters, comprised of the liquid argon LAr barrel electromagnetic calorimeter ($|\eta| < 1.475$) and the Tile hadronic calorimeter ($|\eta| < 1.7$). Three additional calorimeter subsystems are located in the forward regions of the detector: the LAr electromagnetic end-cap calorimeters, the LAr hadronic end-cap calorimeter, and the forward calorimeter comprised of separate electromagnetic and hadronic components.

The precision and accuracy of energy measurements made by the calorimeter system is integral to this analysis and the procedures to establish such measurements are described in Ref. [25]. The baseline electromagnetic (EM) energy scale of the calorimeters derives from the calibration of the electronics signal for the energy deposited by electromagnetic showers. The hadronic calorimeter has been calibrated with electrons and muons in beam tests and the energy scale has been validated using muons produced by cosmic rays with the detector *in situ* in the experimental hall [23]. The invariant mass of the Z boson in $Z \rightarrow ee$ events measured *in situ* in the same data-taking period is used to adjust the calibration for the EM calorimeters.

Dedicated trigger and data acquisition systems are responsible for the online event selection which is performed in three stages: Level 1, Level 2, and the Event Filter. Level 1 utilizes information from the calorimeter and muon systems using hardware-based algorithms. Level 2 and the Event Filter are collectively referred to as the High Level Trigger and utilize software algorithms running on large farms of commercial processors. The measurements presented in this paper rely primarily on the hardware-based Level 1 calorimeter trigger. At this level, coarse calorimeter information is used to reconstruct jets in the trigger system with a square sliding-window algorithm in $\eta - \phi$ space.

4 Data Samples and Event Selection

4.1 Data Sample and Event Selection

The data used for the analysis of event shapes represent the entire 2010 dataset collected at $\sqrt{s} = 7$ TeV and correspond to an integrated luminosity of $\int \mathcal{L} dt = 35.0 \pm 1.1 \text{ pb}^{-1}$ [26].

A sample of events containing high- p_T jets is selected via a Level 1 inclusive single jet trigger with a nominal transverse energy threshold of 95 GeV at the EM scale. The of-

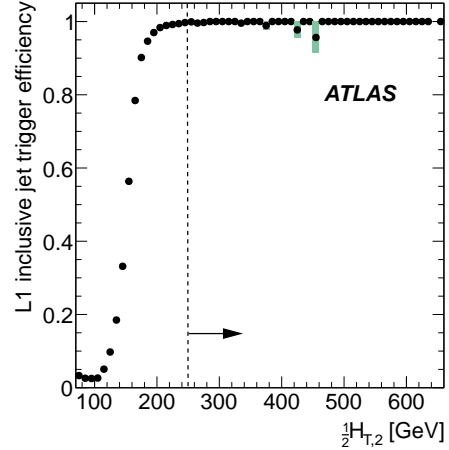


Fig. 1 Inclusive jet trigger efficiency as a function of $\frac{1}{2}H_{T,2}$ evaluated offline. The efficiency to select events with $\frac{1}{2}H_{T,2} > 250$ GeV using this trigger is greater than 99.8%.

fline selection requires two leading jets with a mean transverse momentum $\frac{1}{2}H_{T,2} > 250$ GeV and that the rapidity of each leading jet satisfy $|y| < 1.0$. Subleading jets yield non-zero values of each event shape variable and must have $p_T > 30$ GeV and be within $|y| < 1.5$ in order to be used in the calculations. This choice of event selection is partially driven by the trigger threshold which is fully efficient only at an offline jet transverse momentum of $p_T > 250$ GeV. High momentum jets are also less susceptible to the impact of multiple proton-proton interactions (pile-up) and have a smaller jet energy scale relative uncertainty. The use of $\frac{1}{2}H_{T,2}$ has been shown to be significantly more stable against higher-order corrections to the jet cross section [27]. The inclusive single jet trigger efficiency is evaluated with respect to the offline $\frac{1}{2}H_{T,2}$ selection, as shown in Fig. 1, and is on average greater than 99.8%, thereby removing the need for trigger efficiency corrections. This determination is made *in situ* using a trigger selection with a threshold of 30 GeV at the EM scale.

The presence of pile-up in these data has the potential to impact both event selection and jet reconstruction. Experimentally, reconstruction of primary vertices using tracks measured in the ID provides a measure of the multiplicity of such additional interactions on an event-by-event basis. The 2010 data contain an average number of primary vertices, N_{PV} , of approximately $\langle N_{PV} \rangle = 3$, with a tail extending to $N_{PV} \geq 10$. The vertex with the highest total squared track momentum, $\sum p_{T,track}^2$, is assumed to be the vertex at which the hard scattering that triggered the event occurred.

Two primary effects are expected from pile-up: augmentation of the jet energy scale for jets produced in the hard scattering and pile-up jets produced directly by the additional pp collisions within the same bunch crossing. The consequence of the former is typically an offset to the energy

scale which is corrected as described below. The presence and impact of pile-up jets on the event shape measurements is discussed in more detail in Section 5.2.

4.2 Jet Reconstruction and Calibration

Jets reconstructed with the anti- k_t algorithm [28, 29] are used for the event shape measurements presented here. This algorithm yields regular, approximately circular jets whose boundaries are well described by the nominal jet radius. A jet radius of $R = 0.6$ is used here, as in many Standard Model jet-physics measurements in ATLAS, compared to the smaller radius $R = 0.4$ jets used for a variety of new physics searches and top-quark measurements. This choice is made in order to minimize jet-by-jet corrections due to higher-order emissions and to maximize the reconstruction efficiency since the additional small uncertainties that accompany the larger radius do not significantly impact the measurement. Anti- k_t jets have been shown to be less susceptible than other jet algorithms to systematic effects such as pile-up and close-by jet activity. The inputs to the jet algorithm are topological energy clusters at the EM scale [30]. Following an average offset correction derived *in situ* to account for noise and contributions due to pile-up, an η - and p_T -dependent jet energy correction referred to as the EM+JES correction [25] is applied to all jets to compensate for energy loss in the calorimeters, detector geometry and other effects. Jet quality criteria such as the timing of calorimeter cell signals, the EM energy fraction, and pulse shape information are used to select only those events with well measured jets. These criteria are designed to remove events that are likely to have contamination due to beam-related backgrounds, cosmic rays, or detector defects. In order to further reduce non-collision backgrounds, each event must contain at least one primary vertex consisting of at least five tracks with transverse momenta $p_T^{\text{track}} > 150$ MeV. The MC simulation is reweighted in order to match the primary vertex multiplicity observed in the data.

4.3 Monte Carlo Simulation

Dijet and multi-jet events are generated using two approaches. The first uses direct perturbative calculation of the tree-level matrix elements in powers of the strong coupling constant, α_s . The matrix elements are evaluated at LO in α_s for each relevant partonic subprocess. This is a so-called “multi-leg” method. The second approach implements a sampling of the phase space available for gluon emission with some suitable approximations. The latter uses LO perturbative calculations of matrix elements for $2 \rightarrow 2$ processes and relies on the parton shower implementation to produce the equivalent of multi-parton final states.

The multi-leg technique is used by ALPGEN [31]. In the analysis presented here, ALPGEN 2.13 is used with up to six final-state partons. ALPGEN is interfaced to both HERWIG 6.510 [32] to provide the parton shower and hadronization model, and to JIMMY 4.31 [33] for the underlying event model. The CTEQ6L1 LO [34] parton distribution function (PDF) with LO α_s is used for ALPGEN.

The parton shower simulation programs PYTHIA [35] and HERWIG++ [36] both implement the second approach for QCD jet production and rely on the parton shower to generate multi-jet final states. PYTHIA 6.423 with the Perugia 2010 tune [37] and HERWIG++ 2.4.2 are used to compare to the data; these provide shower models that are p_T ordered and angular ordered, respectively. LO PDFs are taken from the MRST2007 LO* [38, 39] PDF for HERWIG++, and from the CTEQ6L1 LO [34] PDF in PYTHIA.

The MC programs used for comparison to the measurements of event shapes are chosen in part for their ability to describe other ATLAS jet-based measurements. The multi-jet cross section measurements [40], which constitute the same final states as those probed in this analysis, exhibit very good agreement with the predictions from ALPGEN. HERWIG++ not only exhibits good agreement with individual jet shape measurements [41] but is also tuned to yield good agreement with event shape measurements from the Large Electron Positron (LEP) collider experiments. Finally, the Perugia 2010 tune of PYTHIA also shows good agreement with the ATLAS jet shape measurements and has been tuned using the theoretical input from higher-order calculations of event shapes presented in Ref. [11]. These three MC programs thus provide well motivated predictions for the final state observables measured via event shapes.

Events generated by these MC programs are passed through a full simulation [42] of the ATLAS detector and trigger based on GEANT4 [43] and processed in the same way as collision data. The Quark-Gluon String Precompound (QGSP) model [44] is used for high energy inelastic scattering of hadrons by nuclei, and the Bertini cascade model [45] is used to describe the interactions of hadrons with the nuclear medium. Alternative GEANT4 physics lists that specify particle and process definitions, using a combination of the FRITIOF [46] and Bertini models and QGSP without Bertini, are used as part of the studies to understand the uncertainties on the jet energy scale.

5 Results and Systematic Uncertainties

The event shape measurements using jets presented here are corrected to particle-level, after accounting for detector efficiencies and instrumental effects. Particle-level jets are constructed from all final state particles from the MC simulation with lifetimes longer than 10 ps. Direct comparisons can thus be made between the results presented here and

MC generator data after parton shower and hadronization. The dependence of these observables on $\frac{1}{2}H_{T,2}$ is also evaluated. This allows both the investigation of the results and a clearer picture of regions of phase space in the measurement where MC predictions may not fully describe the data.

5.1 Accounting for Detector Effects

In order to compare the predictions of MC event generators with the measurements, several effects must be accounted for. Efficiency loss due to detector coverage and resolution, small detector biases, and other effects, may all affect the measured value of an event shape variable. In order to account for such effects, the MC and detector simulation are used to estimate their impact. MC events after full detector simulation are used to derive bin-by-bin corrections that are applied to the detector-level measurements of each event shape variable to obtain the unfolded, particle-level result to which the MC simulated events after parton shower may be compared. These corrections are typically smaller than 10%. The bin sizes were chosen to be approximately commensurate with resolution, with individual bin purities significantly greater than 50%.

The primary MC generator used for evaluating the corrections is ALPGEN, since the detector-level distributions are well described and ALPGEN models the multi-jet cross-section well [40]. As a cross-check of the method, the corrections evaluated with ALPGEN are compared to those obtained from HERWIG++ and PYTHIA. For $\ln y_{23}$, A , and S , this component of the uncertainty is approximately 2%–8%, which is smaller than both that due to the overall JES systematic uncertainty discussed below and the finite sample size with which the correction factors are determined. However, for the thrust event shape variables, in particular for $T_{m,\perp}$, and S_{\perp} , the generator dependence of these corrections is approximately 10% for the majority of the range of those measurements.

5.2 Systematic Uncertainties

Multiple effects are present in the measurement of event shapes due to the inclusive nature of these observables. These include the uncertainty due to the jet energy scale, the effects of multiple pp interactions, the finite resolution, and the fiducial range of the detector. All of these effects are evaluated and accounted for in the measurement. The dominant uncertainties are the jet energy scale and generator dependence of the corrections in regions of high statistical precision.

The uncertainty on the JES established by the jet calibration procedure [25] influences the final event shape measurement via both the thresholds used to select events and the

momenta used to calculate the event shape observables. This uncertainty is primarily established by the measurement of the single hadron response using test beam data, but is also verified *in situ* during 2010. For jets used in these measurements the typical JES uncertainties are 2.3%–3.0%. The impact of this source of systematic uncertainty is reduced by the explicit use of ratios of jet momenta for several observables, although the jet yield can still vary for a given event due to these selections. Variations of the individual jet momenta are performed within the systematic uncertainties of the JES measurement. For nearly all measured event shape variables, the overall JES uncertainty has the largest impact apart from statistical precision. Most observables have an approximate 5% uncertainty due to the JES, with A and τ_{\perp} being impacted by up to 10% in the steeply falling tails of the distributions.

Additional jets present in the event due to pile-up may also alter the observed event shape. This may be of particular importance for those measurements that are explicitly dependent upon the jet multiplicity, such as those computed from the event transverse momentum tensor. The impact of pile-up on the jet energy is accounted for by the energy scale corrections and uncertainty discussed above, where the pile-up contribution becomes negligible for $p_T > 250$ GeV.

A crucial tool in the identification of jets from pile-up is the jet-vertex fraction, or JVF [47]. This discriminant estimates the contribution of pile-up to a single jet by measuring the fraction of charged particle momentum in the jet that originates in the hard scatter. The rate of additional jets from pile-up in events with between five and eight primary vertices exhibits an increase of a factor of two compared to events with two primary vertices. Because the overall fraction of events with greater than five reconstructed primary vertices is approximately two percent, and these jets tend to have a much softer p_T spectrum, the impact due to additional jets is significantly smaller than other systematic and statistical uncertainties for all measurements.

To further establish the systematic uncertainty incurred by pile-up, comparisons are made between the observed detector-level distributions in events with and without additional reconstructed vertices. A slight variation of a few percent at low $\ln y_{23}$ is observed as well as a relative 10% increase in the fraction of events at higher τ_{\perp} . Similar observations are made by evaluating the impact of the JES uncertainty. Furthermore, each event shape is measured as a function of the JVF of the third jet in the event to directly test the effect of pile-up on the final state observables. Variations of less than 3% are observed when requiring that jets contain a high fraction of associated track momentum originating in the identified hard-scatter vertex. This effect is taken into account in the systematic uncertainty in the final result.

5.3 Event Shape Distributions

The normalized distributions of the third-jet resolution parameter and aplanarity are shown in Figs. 2(a) and (b). In the case of y_{23} , where the primary sensitivity is to the description of the momentum of the third jet, PYTHIA provides the most accurate description of the data, whereas HERWIG++ exhibits slightly better agreement than ALPGEN. Although ALPGEN provides exact tree-level matrix element calculations for up to six jets, it overestimates the fractions of events in the range $\ln y_{23} < -5$. This is qualitatively expected because ALPGEN's more precise calculation of the high jet multiplicity states is primarily concerned with jets near the hard scale of the event. In this region, the impact of the leading-logarithm resummation in PYTHIA and HERWIG++ plays a significant role. As a result, PYTHIA and HERWIG++ both describe the data more accurately than ALPGEN for this event shape variable, particularly at small values of $\ln y_{23}$.

Aplanarity measures the sum of the transverse momenta out of the event plane defined primarily by the two hardest jets. The deviation of the MC prediction from the data is significant for HERWIG++, with some small differences observed with respect to ALPGEN as well. The measurements consistently support more highly aplanar events than predicted by HERWIG++, with the majority of the distribution observed to be significantly different from the MC prediction. The agreement with PYTHIA is good across the full distribution. These results suggest that the event shape is more accurately described by the exact multi-jet prediction provided by the multi-leg matrix element generator (such as ALPGEN) and the model provided by PYTHIA.

The statement that the topological distribution is better described by ALPGEN and PYTHIA than by HERWIG++ is supported by the measurement of the transverse thrust, τ_\perp . Fig. 2(c) exhibits the same behavior as observed in the aplanarity: HERWIG++ predicts fewer than observed highly isotropic events at large τ_\perp . Throughout the distribution, ALPGEN and PYTHIA both predict the measured thrust well. The minor component of the thrust, $T_{m,\perp}$, or the out-of-plane thrust magnitude, shown in Fig. 2(d), does not exhibit as large a difference as observed in the aplanarity. A slight overestimation by PYTHIA is observed for intermediate values of the thrust minor component, $0.25 < T_{m,\perp} < 0.40$.

Lastly, the sphericity and transverse sphericity distributions shown in Figs. 2(e) and (f) exhibit differences between all three generators and the data. The construction of the transverse sphericity as a ratio of eigenvalues of the momentum tensor of the event leads to a slightly improved description. In both cases, PYTHIA provides the best description of the data, and the data are better described by ALPGEN than HERWIG++. In particular, HERWIG++ underestimates the number of highly spherical events in the range $0.40 < S < 0.72$.

5.4 Dependence on $\frac{1}{2}H_{T,2}$

The measurement of the distribution of event shape observables allows for a detailed comparison with MC predictions for a large range of kinematic phase space defined by $\frac{1}{2}H_{T,2}$ in multi-jet events. It is informative to evaluate the explicit dependence of these shapes on the kinematic properties of the event in order to determine potential differences in the modeling of the data for different jet momentum ranges. The evolution of each event shape variable with $\frac{1}{2}H_{T,2}$ exhibits a similar trend as that expected by the running of α_s which leads to a reduction in extra gluon radiation and thus a reduction in the value of each event shape.

Figure 3 depicts the dependence of the mean of each event shape variable on $\frac{1}{2}H_{T,2}$ as it approaches 1 TeV. In all cases, a general trend is observed in which the mean decreases as $\frac{1}{2}H_{T,2}$ increases. This can be understood in terms of how the dynamics of the $2 \rightarrow 2$ process evolves with energy. As the energy in the leading jets increases, the di-jet structure dominates because of kinematics and because α_s decreases as $\frac{1}{2}H_{T,2}$ increases, causing the dominant NLO corrections which generate higher relative momentum gluon emission to decrease as well.

The variation as a function of $\frac{1}{2}H_{T,2}$ is the largest for $\langle \ln y_{23} \rangle$, which indicates a change of y_{23} of nearly a factor of five between $\frac{1}{2}H_{T,2} = 300$ GeV and 800 GeV. Nonetheless, the agreement between the MC prediction and the observed event shape variable dependence is good for all generators. The observation made above that too few highly aplanar events are present in the predictions from HERWIG++ is again observed in the evolution of $\langle A \rangle$ with the momentum scale of the event. The agreement among the mean values measured and the MC predictions improves for $\frac{1}{2}H_{T,2} > 500$ GeV, although the systematic uncertainties are larger and the statistical power of the measurement is reduced. Similarly, the evolution of $\langle \tau_\perp \rangle$ with $\frac{1}{2}H_{T,2}$ in Fig. 3(c) is underestimated by HERWIG++. ALPGEN and PYTHIA consistently predict the mean value of τ_\perp and its evolution with $\frac{1}{2}H_{T,2}$ more accurately, whereas all of the generators describe $\langle T_{m,\perp} \rangle$ vs. $\frac{1}{2}H_{T,2}$ well. Finally, the sphericity and the transverse sphericity (Figs. 3(e) and (f)) are both measured to be approximately 10% larger than predicted by the three MC programs at low $\frac{1}{2}H_{T,2}$, whereas the agreement again improves for higher values. This difference is driven by the underestimate of highly spherical events observed in Figs. 2(e) and (f) which decreases as the average sphericity decreases as a function of $\frac{1}{2}H_{T,2}$.

6 Summary

Six event shape observables are measured with jets in proton-proton collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector with a data sample of 35 pb^{-1} .

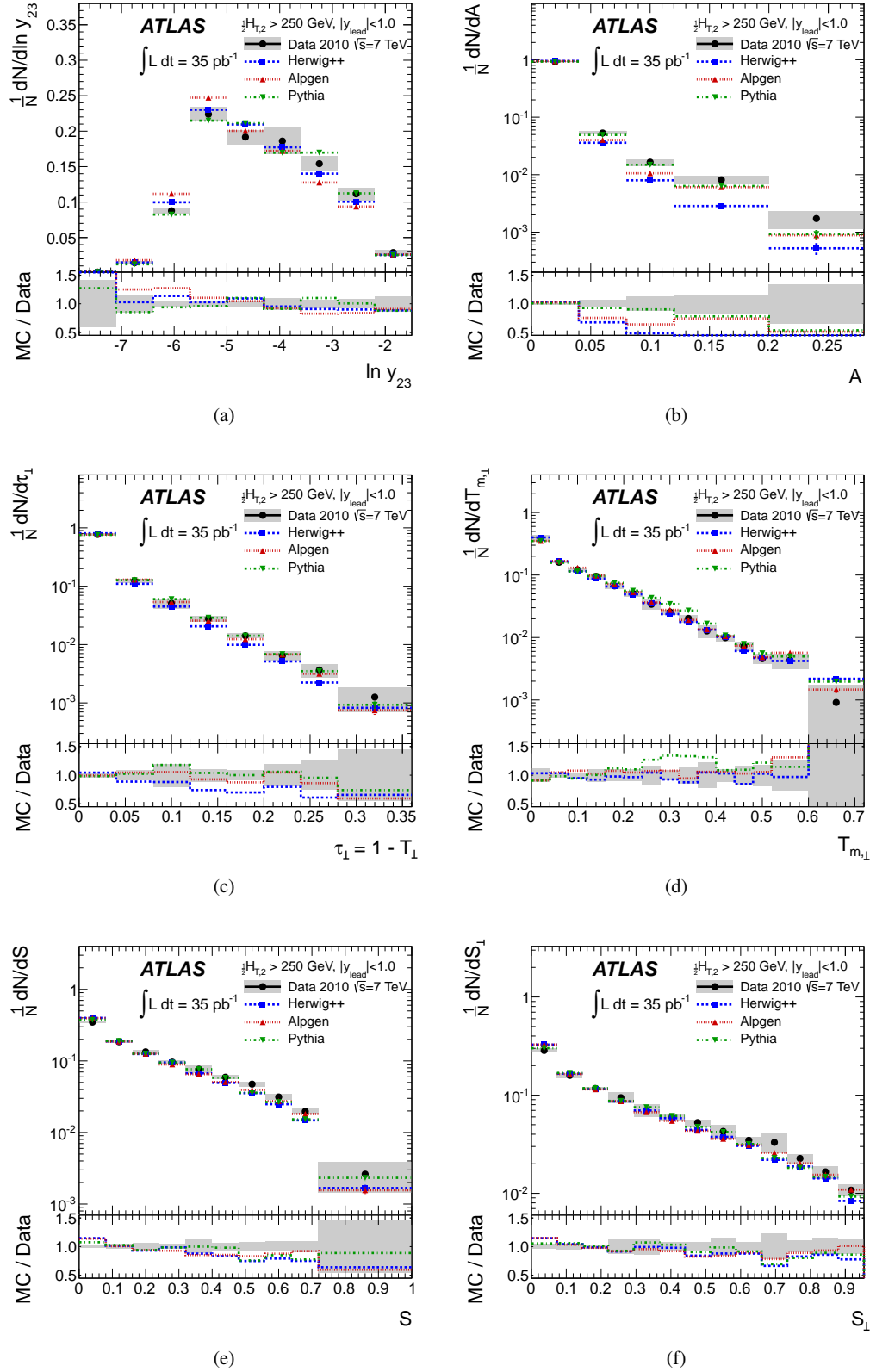


Fig. 2 Unfolded hadron-level distributions of the (a) third-jet resolution parameter, $\ln y_{23}$, (b) aplanarity, A , (c) transverse thrust, τ_{\perp} , (d) minor component of the transverse thrust, $T_{m,\perp}$, (e) sphericity, S , and (f) transverse sphericity, S_{\perp} . The uncertainty shown for the data includes statistical and systematic uncertainties.

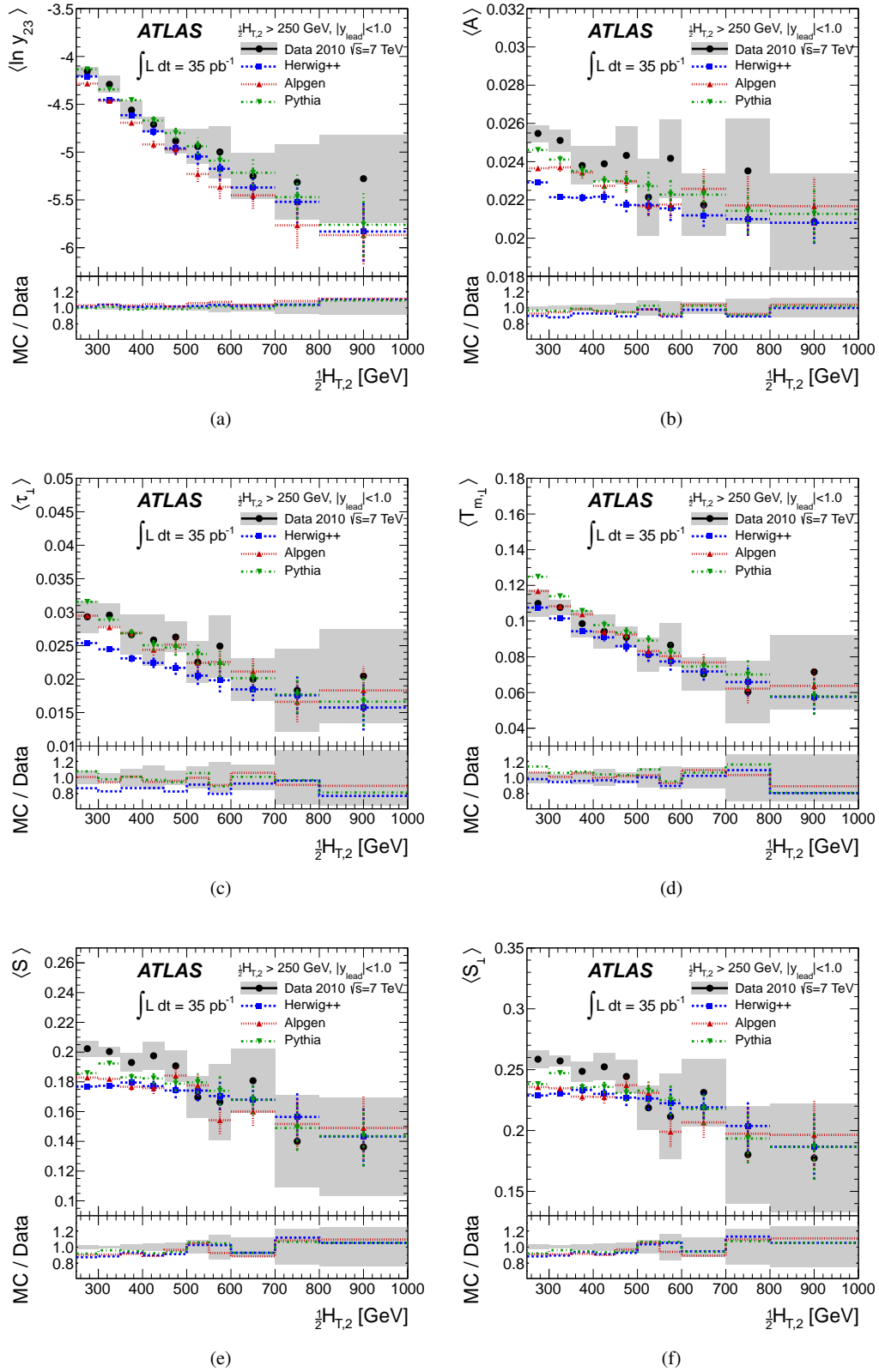


Fig. 3 Mean value of each event shape variable as a function of $\frac{1}{2}H_{T,2}$. Comparisons are made between the MC generators HERWIG++, ALPGEN and PYTHIA.

Measurements are performed up to an event $H_{T,2}$ of 2 TeV and are compared with different Monte Carlo event generators. Overall shape comparisons are made with these MC programs, as well as the kinematic evolution of the mean value of each event shape variable with $\frac{1}{2}H_{T,2}$. Reasonable agreement is observed in most kinematic and topological regions. The measurements suggest that the modeling of the data by PYTHIA (Perugia 2010) and ALPGEN are more accurate than that by HERWIG++, in particular for the planarity, A , and transverse thrust, τ_{\perp} . The good description provided by ALPGEN (+HERWIG/JIMMY) of the multi-jet cross-section [40] is reflected as well in these measurements, although the description provided by PYTHIA tends to model the data more accurately. PYTHIA predicts a slightly higher mean $T_{m,\perp}$ at low $\frac{1}{2}H_{T,2}$ than observed in the data whereas HERWIG++ predicts a slightly lower mean thrust. The systematic uncertainties in the measurement of S and S_{\perp} are found to be relatively small. The observation that the measured mean value of each event shape variable decreases with $\frac{1}{2}H_{T,2}$ is consistent with the trend expected from the running of α_s and is generally well modeled by the MC simulations.

Comparisons of these results to previously published LHC measurements of hadronic event shapes [7] indicate that the slight overestimate by PYTHIA of events with thrust minor component, $T_{m,\perp}$, in the range $0.25 < T_{m,\perp} < 0.40$ is observed in each case. However, the good agreement observed in this study between data and both PYTHIA and ALPGEN for the thrust, τ_{\perp} is not seen in Ref. [7]. The different tunes of HERWIG++ and PYTHIA, as well as the different underlying event and hadronization models interfaced to ALPGEN in the two measurements, may account for these differences. Furthermore, the systematic uncertainties associated with these measurements are in many cases similar to the differences between the generators.

ATLAS measurements of the jet shape [41], jet fragmentation function [48], and multi-jet cross-section also provide additional insight into the results shown here. PYTHIA and HERWIG++ 2.4.2 both provide a reasonable description of the fragmentation function of high momentum jets in the data, whereas only the former models the jet shape accurately. ALPGEN, on the other hand, yields a fairly accurate description of the three-jet to two-jet cross-section ratio [40], although it produces jet shapes that are significantly narrower than those measured here.

These results, supported by other ATLAS measurements of the hadronic final state, reinforce the importance of the leading-logarithm resummation in parton shower MC event generators. They also demonstrate the ability of leading order MC to provide a reasonable description of multi-jet event shapes.

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